

Exchange and
CPT

Daniel L. Miller

Compare
 $SO(3N, N)$ with
 $N \times SO(3, 1)$

Embedding

CPT

Exchange

Commutations
between CPT
and Exchange

$SO(3N, N)$

$N \times SO(3, 1)$

Conclusions

matter-antimatter
asymmetry

Call for experiments

Few notes

Proof of anticommutation between exchange and charge conjugation for spinors

Interactions between identical Particles (Matter) and
between identical Anti-Particles (Anti-Matter) are not the
same.

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Outline

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This work is based on various group embeddings published in 1990–2010 by prof. M. Berry and other teams, and author works from 2014-2016; preprint was 1st submitted in 2019

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Few notes

- 1 Compare $SO(3N,N)$ with $N \times SO(3,1)$
 - Embedding
 - CPT
 - Exchange
- 2 Commutations between CPT and Exchange
 - $SO(3N,N)$
 - $N \times SO(3,1)$
- 3 Conclusions
 - matter-antimatter asymmetry
 - Call for experiments
 - Few notes

Compare $SO(3N,N)$ with $N \times SO(3,1)$ – Embedding

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$N \times SO(3,1)$ world can be embedded into $SO(3N,N)$; It gives rise to spinor branching rules; we will not need them explicitly; just compare γ -matrices.

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$SO(3N,N)$

Bispinor has 2^{2N} components
Vector has $4N$ components

$$\gamma_j^\mu = \underbrace{\gamma^{\text{FIVE}} \otimes \dots \otimes \gamma^\mu}_{j-1 \text{ times}} \underbrace{\dots \otimes \hat{1}}_{N-j \text{ times}}$$

$$\{\gamma_j^\mu, \gamma_k^\nu\} = 2\delta_{jk}\delta_{\mu\nu}$$

γ -matrices are different in their commutation relations, but dimensions and amounts are the same.

$N \times SO(3,1)$

$N \times$ Bispinor has 2^{2N} components
 N Vectors have $4N$ components

$$\gamma_j^\mu = \underbrace{\hat{1} \otimes \dots \otimes \gamma^\mu}_{j-1 \text{ times}} \underbrace{\dots \otimes \hat{1}}_{N-j \text{ times}}$$

$$[\gamma_j^\mu, \gamma_k^\nu] = 0, j \neq k$$

$$\{\gamma_j^\mu, \gamma_j^\nu\} = 2\delta_{\mu\nu}$$

Compare $SO(3N,N)$ with $N \times SO(3,1)$ – CPT

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CPT transformations are defined in a similar way

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$$\begin{aligned}\hat{C} &= \prod_j \hat{C}_j, \quad \hat{C}_j \sim \gamma_j^2 & \hat{P} &= \prod_j \hat{P}_j, \quad \hat{P}_j \sim \gamma_j^0 & (1) \\ \hat{T} &= \prod_j \hat{T}_j, \quad \hat{T}_j \sim \gamma_j^1 \gamma_j^3 & \hat{I}^4 &= \prod_j \hat{I}_j^4, \quad \hat{I}_j^4 \sim \gamma_j^0 \gamma_j^1 \gamma_j^2 \gamma_j^3 \\ & & \hat{I}^4 &= \hat{C} \hat{P} \hat{T}\end{aligned}$$

SO(3N,N) vs $N \times \text{SO}(3,1)$ – Exchange

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Exchange is $2^{2N} \times 2^{2N}$ matrix operator acting on bi-spinor and its coordinates.

It exchanges components of coordinates of these bi-spinors in the same way for SO(3N,N) and $N \times \text{SO}(3,1)$

$$\hat{E}_{ij} : \begin{matrix} x_j^\mu \rightarrow x_i^\mu, & x_i^\mu \rightarrow x_j^\mu, \\ x_k^\mu \rightarrow x_k^\mu, & k \neq i, j \end{matrix} \quad (2)$$

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SO(3N,N)

The Exchange operator can be derived explicitly; properties

$$\hat{E}_{ij}^x \gamma_j^\mu \hat{E}_{ij}^x = \gamma_i^\mu \quad (3)$$

$$\hat{E}_{ij}^x = \hat{E}_{ji}^x \quad \hat{E}_{ij}^x \hat{E}_{ij}^x = \hat{1}$$

$N \times \text{SO}(3,1)$

The Exchange operator E^h was postulated in 1926 by Werner Heisenberg for non-relativistic quantum mechanics. It can be generalized for QED



Commutations between C, P, T, CPT and Exchange

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The anticommutation of the γ -matrices between themselves, along with Eqs. (1,3), gives for SO(3N,N)

$$\begin{aligned}\hat{E}^x \hat{C} \hat{E}^x &= -\hat{C} & \hat{E}^x \hat{P} \hat{E}^x &= -\hat{P} \\ \hat{E}^x \hat{T} \hat{E}^x &= \hat{T} & \hat{E}^x \hat{\gamma}^4 \hat{E}^x &= \hat{\gamma}^4\end{aligned}\quad (4)$$

For N x SO(3,1) Exchange commutes with all CPT symmetries, this is must for proof of the the spin-statistics theorem.

$$\begin{aligned}\hat{E}^h \hat{C} \hat{E}^h &= \hat{C} & \hat{E}^h \hat{P} \hat{E}^h &= \hat{P} \\ \hat{E}^h \hat{T} \hat{E}^h &= \hat{T} & \hat{E}^h \hat{\gamma}^4 \hat{E}^h &= \hat{\gamma}^4\end{aligned}\quad (5)$$

Symmetry lowering $SO(3N,N) \rightarrow N \times SO(3,1)$ and the proof of anticommutation between exchange and charge conjugation

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Central hypothesis

- Let's postulate our world was initially $4N$ -dimensional, $SO(3N,N)$
- Then the Big Bang reduced the spatial symmetry to $SO(3,1)$, giving rise to N spinor fields, namely $SO(3,1)$ bi-spinors
- Let's assume the Exchange operator \hat{E} and all many body CPT transformations are preserved upon the symmetry lowering

$$E = E^X \Rightarrow \hat{E}\hat{C}\hat{E} = -\hat{C}, \quad \{\hat{E}, \hat{C}\} = 0$$

- This ends the proof of anticommutation between exchange and charge conjugation

Consequences for matter-antimatter asymmetry

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$$\{\hat{E}, \hat{C}\} = 0 \quad \{\hat{E}, \hat{P}\} = 0 \quad [\hat{E}, \hat{T}] = 0$$

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- the charge conjugation reverses the statistics (for spinors)
- bosonic fields annihilate fermionic fields
- no changes in single line diagrams (eg magnetic moment)
- no Periodic Table for antimatter; collapse of antiatom size $Z^{-1/3} \rightarrow Z^{-1}$ due to lack of the degeneracy pressure
- this explains matter-antimatter asymmetry within QED
- it would be no selfconjugated (truly neutral) spinors
- vector fields stay bosonic (contract matter with matter)
- scalar contracting matter with antimatter must be the supercharge

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Few experiments can either confirm or invalidate anticommutation between exchange and charge conjugation. As of today none of such evidences is available.

Possible evidences:

- spectroscopy of antihydrogen molecule
- positron-positron Møller scattering
- compare internal degeneracy of proton and anti-proton

Lorentz invariance

The Exchange operator derived by making use of rotations in $SO(3N,N)$ is Lorentz invariant in $SO(3,1)$

Many body CPT

Many body CPT transformations Eq. (1) as they defined for $SO(3N,N)$ can be used for $N \times SO(3,1)$. So the assumption that CPT transformations are preserved upon the symmetry lowering is merely semantical. The assumption that the Exchange is preserved upon the symmetry lowering is critical.

